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# Quantum chemical origin of high ionization potential and low electron affinity of Tungsten Hexafluoride

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# ABSTRACT

This paper deals with quantum chemical explanation of high ionization potential and low electron affinity of  $WF_6$  as compared other species of this family. Calculated structure and vibrational frequencies of  $WF_6$  by B3LYP, B3PW91 and MP2 schemes are compared with experimental values which establish superiority of B3PW91 over B3LYP and MP2 in the present case. Anionic and cationic forms of  $WF_6$  are optimized and analysed at B3PW91 level. It is revealed that low electron affinity of  $WF_6$  is due to extra electron localization on W which goes to antibonding molecular orbital. It is also shown that highest occupied molecular orbital is delocalized over F atoms and it is not easy to extract electron from electronegative F which results in very high ionization potential of  $WF_6$ .

Keywords: Tungsten hexafluoride, Electron affinity, Ionization Potential, Molecular Orbital, Atomic charge distribution.

# INTRODUCTION

In the periodic table, transition metal elements have a great importance due to their variable coordination number. They form a variety of compounds which are referred to as coordination complexes. The nature of bonding in coordination complexes may vary depending upon the nature of its components, nature of metal and their oxidation states and ligands. Transition metal hexafluorides constitute an important class of coordination complexes, a central metal element is surrounded by six F ligands. The most common feature of hexafluorides includes an octahedral structure with high electron affinity (EA). However, octahedral symmetry may be broken due to an effect, so called, Jahn-Teller distortion, the high electron affinity of such species still gets attention of a wide community. Many transition metals have been reported to form hexafluoride molecules possessing very high electron affinities as compared to fluorine, at least, theoretically [1-7]. For instance, we have reported a systematic study on Osmium (Os) with a number of F ligands and found that OsF<sub>6</sub> is distorted to square planar ( $D_{4h}$ ) geometry with adiabatic EA of about 6.5 eV [8]. Recently, our study [9] on WF<sub>n</sub> (n = 1-5) species indicate that the tungsten fluorides behave differently than other transition metal fluorides.

Tungsten hexafluoride (WF<sub>6</sub>) is well known for its distinguished properties. WF<sub>6</sub> is widely employed in the manufacture of semiconductor devices. It is used for deposition of tungsten metal to form interconnects due to its high conductivity and compatibility with the silicon substrate. In contrast to other hexafluorides of the same series, it is perfect octahedron. Moreover, it possesses very low electron affinity and very high ionization potential [1]. This is indeed a strong indication of its enormous stability against addition or removal of an electron from WF<sub>6</sub>. In the present work, we perform a quantum chemical survey on these distinguished electronic properties of WF<sub>6</sub>. We have tested various methods (B3LYP, B3PW91 and MP2) in order to choose appropriate scheme for the problem under study. To benchmark our calculations, we have compared calculated geometry and vibrational frequencies of WF<sub>6</sub> with corresponding experimental values. Finally, using atomic charge distribution and molecular orbital analysis at

the most appropriate level of theory, we have explained the origin of high ionization potential and low electron affinity of  $WF_6$ .

# **COMPUTATIONAL DETAILS**

All computations were carried out with Gaussian 09 program [10] within ab initio density functional schemes. The functionals, in which Becke's three parameter exchange term (B3) [11] is combined with correlation terms of Lee, Yang and Parr (LYP) [12] as well as of Perdew and Wang (PW91) [13], were used. In addition, we have applied second order perturbation method of Moller and Plesset (MP2) [14]. All atom basis set of Stuttgart-Dresden-Dunning (SDD) was employed throughout these calculations. The visual animations of vibrational modes and all relevant graphics were generated with the help of GaussView 5 [15] and Chemcraft 1.8 packages [16].





Figure 1. Model structure of Tungsten hexafluoride (WF<sub>6</sub>)

## Molecular structure and vibrations

The molecular structure of WF<sub>6</sub> is optimized by B3LYP, B3PW91 and MP2(FC) methods using SDD basis set. All methods lead to an octahedral ( $O_h$ ) structure of WF<sub>6</sub> with singlet spin state as shown in Fig. 1. The calculated bond-lengths are 1.884, 1.878 and 1.903 A° by B3LYP, B3PW91 and MP2 methods, respectively. The corresponding experimental value is 1.833 A° measured by electron diffraction method [17]. Thus, B3PW91 calculated bond-length is more close to experimental value and highly overestimated by MP2 scheme.

Vibrational frequency calculations are also performed at the same level of theories within the harmonic approximation. Calculated harmonic frequencies and intensities are listed in Table 1. For a comparison, corresponding experimental values [18] are also included. The transition corresponding to symmetric stretching  $(a_{1g})$  is calculated at 716 cm<sup>-1</sup> by B3LYP, 722 cm<sup>-1</sup> by B3PW91 and 699 cm<sup>-1</sup> by MP2 against experimental value of 771 cm<sup>-1</sup>. Very strong mode  $(t_{1u})$  observed at 712 cm<sup>-1</sup> is triply degenerate due to asymmetric stretching along three axes. A doubly degenerate mode  $(e_g)$  observed at 677 cm<sup>-1</sup> also corresponds to asymmetric stretching, but along two planes.

Table 1. Vibrational modes of WF <sub>6</sub> calculated at B3LYP, B3PW91 and MP2 level with SDD basis set. Experimental values	are also
included for comparison	

Modes	B3LYP		B3LYP B3PW91		MP2		Experiment <sup>c</sup>	
	Freq.ª	Intensity <sup>b</sup>	Freq.	Intensity	Freq.	Intensity	Freq.	Intensity
alg	716	0.0	722	0.0	699	0.0	771	inactive
tıu	695	222.6	700	226.7	682	260.7	712	v. strong
eg	645	0.0	651	0.0	629	0.0	677	inactive
t <sub>2</sub> g	288	0.0	291	0.0	291	0.0	320	inactive
$t_1 u$	223	39.7	225	39.4	231	39.3	258	strong
t <sub>2</sub> u	115	0.0	116	0.0	123	0.0	127	inactive

<sup>a</sup>Frequency in cm<sup>-1</sup>; <sup>b</sup>Intensity in a.u.; <sup>c</sup>Ref. [18].

All other lower modes associated with bending and distortions are triply degenerate due to symmetry of vibrations. Also due to highly symmetric structure of  $WF_6$ , most of IR modes become inactive. From Table 1, one can see that for  $a_{1g}$ ,  $e_g$ ,  $t_{2g}$  and  $t_{2u}$  modes, IR intensities vanish. In general, our computational methods reproduce and assign experimental bands correctly. The differences between calculated and observed frequencies are mainly due to anharmonicity of vibrations, which go on decreasing for lower frequencies for all three methods employed. It should be noticed, however, that B3PW91 calculated frequencies are more close to experimental values as compared to B3LYP and MP2.

#### Ionization potential and electron affinity

In order to discuss the effect of addition and removal of an electron, we have optimized the structures of anionic and cationic  $WF_6$  at B3PW91 level of theory. The equilibrium geometries of  $WF_6^-$  and  $WF_6^+$  are displayed along with neutral  $WF_6$  in Fig. 2. Evidently, addition of an electron does not affect the molecular geometry of  $WF_6$ . It results only in a marginal increase in the bond-length (by 0.01 A°). On the other hand, removal of an electron leads to change the structure into complex form, ( $WF_4$ )F<sub>2</sub>. In this complex, two F atoms go away from the central W forming a  $F_2$  moiety with the distance of 1.935 A°. The distance of this moiety from centre is increased to 2.302 A°, whereas other bond-lengths are decreased to 1.843-1.851 A°.



Figure 2. Equilibrium geometries of WF<sub>6</sub>, WF<sub>6</sub><sup>-</sup> and WF<sub>6</sub><sup>+</sup> at B3PW91/SDD level. Atoms are labeled with NBO charges (in e) and bondlengths (in A<sup>o</sup>) are also shown

We discuss the origin of high ionization potential and low electron affinity of WF<sub>6</sub>. We analyze the distribution of extra electron over W and F atoms referring to charges labeled on atoms in Fig. 2. These atomic charges are calculated by natural bonding orbital (NBO) analysis [19,20] which are more reliable due to their less basis set dependency. In neutral WF<sub>6</sub>, atomic charge on W is +2.08 e and -0.35 e is contained by each F atom. First we consider the case of removal of an electron, i.e. WF<sub>6</sub><sup>+</sup>. The ionization potential (IP) of WF<sub>6</sub> is calculated to be about 15 eV which is consistent with the experimental value of  $15.24 \pm 0.10$  eV by Hildenbrand [21]. In order to explain such a high IP, we examine the charge distribution upon ionization of WF<sub>6</sub>. The electron is mainly contributed by 2p atomic orbital of two F atoms which forms a F<sub>2</sub> moiety. W atom contributes only 0.1 e during the process of ionization which results in very high IP of WF<sub>6</sub>.



Figure 3. Molecular orbital diagrams of WF<sub>6</sub> and WF<sub>6</sub><sup>-</sup> at B3PW91/SDD level

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Similarly, the EA of WF<sub>6</sub> has been estimated as 3.9 eV which is slightly higher than experimentally estimation of 3.4 - 3.7 eV [22-24]. However, note that this value is very small as compared to those of other hexafluorides of 5d series. For instance, the EA of PtF<sub>6</sub>, 6.85 eV [25] is almost two times that of WF<sub>6</sub>. In general, large EA of hexafluorides results due to extra electron delocalization over several F atoms, leaving electron deficiency at the centre. But this is not the case for WF<sub>6</sub> in which extra electron is completely localized on W atom. The NBO charge difference on W in WF<sub>6</sub> and WF<sub>6</sub><sup>-</sup> is 0.9 e (see Fig. 2). Thus 90% of extra electron is located on W and only 10% is shared equally by peripheral F atoms. The extra electron localization on W disables it to be electron deficient, consequently lowering the EA of WF<sub>6</sub>.



Figure 4. HOMO (left) and LUMO (right) plots of WF<sub>6</sub> and WF<sub>6</sub><sup>-</sup> calculated at B3PW91/SDD level

Let us now discuss these electronic properties from molecular orbital (MO) perspectives. W possesses an outer shell configuration of  $5d^4 6s^2$ . Due to sextet ( $5d^56s^1$ ) configuration as its ground state, it can easily bind with six F ( $2s^22p^4$ ) atoms. In WF<sub>6</sub>, due to the interaction of d orbitals of W and p orbitals of F, d energy levels split into two sets. First set  $e_{1g} (d_x^2 \cdot y^2, d_z^2)$  is 0.64 eV above the second  $t_{1g} (d_{xy}, d_{yz}, d_{xz})$ . The MOs are constructed from metal and ligand orbitals symmetry as shown in Fig. 3 for WF<sub>6</sub> and WF<sub>6</sub><sup>-</sup>. We focus on frontier MOs as they are mainly responsible for chemical reactions or interactions. The highest occupied molecular orbital (HOMO) represents ability to donate an electron while lowest unoccupied molecular orbital (LUMO) denotes ability to accept it. The energy gap between HOMO and LUMO can be used as a stability index of molecule. Figure 4 displays the HOMO-LUMO surfaces for WF<sub>6</sub> and WF<sub>6</sub><sup>-</sup>. Apparently, the HOMO of WF<sub>6</sub> is delocalized over F atoms i.e. ligand MO. Due to strong electronegative nature of F atoms, it is not favorable to extract an electron which is the reason behind high IP of WF<sub>6</sub>.

Addition of an electron to  $WF_6$  causes to shift  $e_{1g}$  and  $t_{1g}$  levels increasing their separation to 0.78 eV, consequently, increasing the energies of HOMO and LUMO. This extra electron goes to antibonding MO of  $WF_6$  as shown in Fig. 3 and in the HOMO plot of  $WF_6^-$  as well (see Fig. 4). This causes to slightly destabilize the molecule due to repulsion created by antibonding electron which is reflected by decrease in HOMO-LUMO gap of  $WF_6$  by addition of an electron. Note the HOMO-LUMO gaps of  $WF_6$  and  $WF_6^-$  are 6.53 eV and 2.18 eV, respectively. This also explains the low EA value of  $WF_6$ .

#### CONCLUSION

In summary, we have used B3PW91 method to explain the high ionization potential and low electron affinity of WF<sub>6</sub>. The reliability of B3PW91 over B3LYP and MP2 is established by calculating structure and vibrational

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frequencies of  $WF_6$ . We have shown that low electron affinity of  $WF_6$  is due to extra electron localization on W and highest occupied molecular orbital is delocalized over F atoms which results in very high ionization potential of  $WF_6$ . Thus present study provides a theoretical explanation of distinguished electronic properties of tungsten hexafluoride.

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## **CONFLICT OF INTEREST**

The authors declare that there is no conflict of interests regarding the publication of this paper. Also, they declare that this paper or part of it has not been published elsewhere.

## **CONTRIBUTION OF THE AUTHORS**

AKS performed all calculations and prepared the draft. AKP contributed in writing the paper. NM designed the research and finalized the draft. AKS, AKP and NM revised the draft paper. All authors read and approved the final version.

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